

Mars and Venus Entry Simulation Capabilities of IRS Plasma Wind Tunnel PWK3

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Abstract

An assessment is made for the inductively driven plasma wind tunnel PWK3 with the goal to derive relevant mass specific enthalpies for typical Mars and Venus atmospheric entry missions. For this purpose an integral method has been used which links the plasma power to the radial distribution of total pressure and fully catalytic heat flux in the plasma jet on basis of a relation from Marvin and Pope. Rebuilding the enthalpies with this relation allows for the derivation of a gas specific proportionality factor. This factor enables the derivation of the mass specific enthalpies at the centre line and the radial profiles for the respective condition are not necessarily required any more. Correspondingly a review of reference CO₂ plasma conditions obtained in past investigations at IRS leads to the identification of an operational envelope in terms of the mass specific enthalpies which are from an energy consideration the prerequisite for the creation of similarities with respect to the real atmospheric entry maneuvers. The analysis shows that PWK3 is capable to cover the full range of mass specific enthalpies that are required for typical Mars and Venus atmospheric entry scenarios.

Keywords: Atmospheric entry simulation, Non-equilibrium plasma, Simulation facilities

1. Introduction

Electrodeless inductively heated plasma generators enable basic thermal protection system material tests (e.g. catalysis- see Herdrich, Fertig, Lein, Steinbeck, Pidan (2010) and Herdrich, Fertig, Petkow, Steinbeck, (2010)) and the simulation of atmospheres that contain chemically reactive components such as of Mars or Venus with a high amount of CO₂ in the atmosphere. The inductively driven generators at IRS have an optimized design where the induction coil is closer to the plasma than it is with other designs, see Herdrich, Petkow (2008) and Nawaz, Herdrich (2009). Therefore, the electromagnetic field loss is reduced. The water cooling system surrounds both the coil and the plasma tube.

In the future, several planetary probe missions, soft landing missions and sample return missions such as "Venus Sample Return Mission", "Mars Mini-Probes", Mars Society Balloon Mission or "Mars Sample Return Mission" (MSR) will be accomplished, see Herdrich, Auweter-Kurtz, Endlich (2003). The current reference mission of ESA, ExoMars, has the final goal of a sample return. The sample, however, has to be brought from the Martian surface such that the overall mission has to cope with both an atmospheric entry at Mars and a hyperbolic re-entry for Earth as e.g. outlined by Bouilly, Bonnefond, Boulier, Balembois, Boquet, Plaindoux, Mignot, (2009).

For such missions both TPS and environment (plasma) during the entry have to be investigated by means of computational and ground facility based simulations. Such ground facilities are the IRS plasma wind tunnels PWK 1-4 reproducing the thermal, aerodynamic and chemical load on the surface of a space vehicle entering an atmosphere. They are operated with different plasma generators; see Herdrich, Auweter-Kurtz, Endlich, (2003) and Auweter-Kurtz, Herdrich, Laure, Wagner (2004).

In addition, plasma wind tunnels can be applied for the development of in-flight instruments e.g. aiming for the

assessment of the conditions at the hot structures and the plasma conditions experienced by the vehicle during entry as e.g. described by Auweter-Kurtz, Fertig, Herdrich, Laux, Schöttle, Wegmann, Winter (2003), Lein, Reimer, Stubicar, Deuble, Auweter-Kurtz, Herdrich, Winter (2009) and Herdrich, Fertig, Lein, Löhle, Preci, Steinbeck, Wernitz, Auweter-Kurtz, Roeser (2010). Furthermore, they can be also used to support development activities for instruments to be used for airborne observation campaigns, see Herdrich, Fertig, Lein, Löhle, Preci, Steinbeck, Wernitz, Auweter-Kurtz, Roeser (2010). Reference conditions can be used for validation of numerical tools as well; see Fertig, Herdrich (2009).

Non-intrusive measurement techniques like emission spectroscopy, Fabry-Perot interferometry and laser-induced fluorescence are used to investigate the plasma flows. Relevant techniques as such are e.g. described and applied by Löhle, Eichhorn, Steinbeck, Lein, Herdrich, Röser, Auweter-Kurtz (2008) and by Matsui, Takayanagi, Komurasaki, Arakawa, Knapp, Herdrich, Auweter-Kurtz (2008). They are applied to determine atomic and molecular density and the velocity distribution. The laser absorption spectroscopy technique of the Department of Aeronautics and Astronautics (Tokyo University) was used for IRS-PWK3 to determine densities and translational temperatures e.g. of atomic oxygen, see also Matsui, Komurasaki, Herdrich, Auweter-Kurtz (2005).

Besides the non-intrusive measurement techniques, mass spectrometry, electrostatic and radiation probes belong to the group of intrusive measurement techniques. Mechanical probes are among the most important instruments for plasma-diagnostic measurements and are often used. Besides the standard sample support system which carries the TPS material sample to be tested, probes for Pitot pressure, Mach number, heat flux, enthalpy and oxygen partial pressure determination are used. Electrostatic probes are used to ascertain the plasma potential, electron density and temperature, energy distribution of the electrons, ion temperature and flow velocity. Radiometric probes are unavoidable when the radiation heat flux cannot be neglected compared to the convective part. This is the case when during sample return missions the entry speed into the Earth's atmosphere is especially high or when the atmosphere of another celestial body (that is to be entered) contains strong radiating species, as for example the atmosphere of Titan.

2. Experimental Set-up

2.1 Facility PWK3

An extensive description of the facility has been developed by Herdrich (2004).

The PWK3 set-up shown in Fig. 1 consists of the IPG plasma source and the vacuum chamber. The chamber is about 2 m in length and 1.6 m in diameter. Optical accesses enable the investigation of the plasma. A heat exchanger between the chamber and the vacuum system cools down the hot plasma to protect the vacuum system from being damaged. The flat lid of PWK3 (left side chamber) is equipped with the IPG and the external resonant circuit consisting of the capacitors with the connection to the IPG coil. The right side flange of the vacuum chamber is connected to the IRS vacuum pump system simulating pressures at altitudes up to 90 km (Earth entry). Total suction power of the pumps amounts to $6 \cdot 10^3$ m³/h at atmospheric pressure and reaches about $2.5 \cdot 10^5$ m³/h at 10 Pa measured at the intake pipe of the system, which has a diameter of 1 m. The base pressure is about 0.5 Pa. The desired tank pressure can be adjusted between the best achievable vacuum and 100 kPa by removing one or more pumps from the circuit and/or mixing additional air into the system close to the pumps.

The external resonant circuit is water cooled. With this, the capacitors, which have a capacity of $6 \text{ nF} \pm 20\%$ each, and the coil are cooled. The resonant circuit depicted in Fig. 2 is built in Meissner type switching using a metal-ceramic triode with an oscillator efficiency of about 75%. Its nominal frequency can be changed by switching the number of capacitors k as well as by the use of coils with different inductivities, i.e. variation of coil turns, n .

For the present investigations, a water cooled 2.5-turn ($n = 2.5$) and 5.5-turn ($n = 5.5$) coil, each having a length l of 120 mm, were used. These configurations lead to inductivities L of 0.45 and 1.8 μH , respectively. The whole circuit is connected to a 375 kW power supply. The incoming anode power can be adjusted by the control of the anode (plate) voltage. Correspondingly, the frequency can be approximated by

$$\omega = 2\pi f \approx \frac{1}{\sqrt{LC}} \quad (1)$$

Here, ω is the angular frequency while f is the frequency. The coil inductivity is calculated by

$$L_{coil} = d \alpha \left(\frac{d}{l} \right) n^2 \cdot 10^{-7} = g(d, l) n^2 \quad (2)$$

as the well-known relationship for long coils is not applicable anymore, see also reference Herdrich, Auweter-Kurtz, Endlich (2003). The coil's diameter d has to be inserted in cm leading to the dimension Henry for L . The parameter α , a correction factor, has a magnitude of 6, as e.g. derived by Herdrich, Petkow (2008). Introducing the parameter k for the number of capacitors used during operation, the frequency can be determined by

$$f_{n,k} = f(n, k) = f_{1,1} \frac{1}{n\sqrt{k}} \propto \frac{1}{n\sqrt{k}} \quad (3)$$

Usually, there are peripheral inductivities due to the coil's input leads which are not negligible. Therefore, it is useful to use Eq. (3) in empirical form depending on k i.e. for coils with specific n .

With the coils used in this investigation ($n = 2.5$, $n = 5.5$), Eqs. (1) and (2) deliver frequencies that are higher than the measured ones. This can be explained by the presence of peripheral inductivities which can be of the same order of magnitude as the coil inductivities:

$$L_{tot} = L_{coil} + L_{Peri} \quad (4)$$

Eqs. (2-4) in combination with the measured frequencies for $n = 2.5$ and $n = 5.5$ lead to the frequency range shown in Fig. 3 for PWK3 using the inductively heated plasma generators IPG3, IPG4 and IPG5.

Figure 3 shows the measured frequencies of the facility PWK3 together with the calculated frequencies using Eq. (1) for different plasmas. Here, Eq. (4) was applied and L_{Peri} was calculated for all measured values. These values depend on the number of capacitors k and on the number of coil turns n . According to Fig. 3, this leads to a frequency range between 0.5 and 1.4 MHz. The usage of just one capacitor together with a 2.5-turn coil ($L_{coil} \approx 0,4 \mu\text{H}$) makes a frequency of 1.9 MHz achievable. This value cannot be operated permanently due to limitations of the energy supply system and current limiting effects that are ongoing with the capacity reduction. The second coil's ($n = 5.5$) inductivity is $L_{coil} \approx 1,8 \mu\text{H}$. The frequencies shown with filled symbols were measured using a looped conductor connected to an oscilloscope. The open square symbol represents frequencies measured with a modified Pearson current monitor described by Herdrich (2004).

Here, the coil currents for $k = 4$ were measured with high time resolution for oxygen and nitrogen while varying the mass flow rate and anode power such that frequency spectra could be obtained by Fourier analysis. For $k = 7$ a Rogowski coil was used to determine the frequencies for argon and air while varying the plate power. Regression analysis was performed for the data by Herdrich (2004). With $L_{peri} \approx 0.75 \mu\text{H}$ Eqs. (1-4) correspond well to the measured frequencies.

Inductively heated plasma generators basically consist of a coil surrounding a plasma container (tube) and capacitors. The alternating current in the coil induces a mostly azimuthal electric field inside the tube. This electric field initiates an electric discharge in the gas that is injected on one side into the tube (Fig. 4). The produced plasma is expanded into the vacuum chamber. The plasma current amplitude - and thus the Ohmic heating - strongly depends on the electric conductivity of the plasma and the resonant frequency.

The principal parts of the plasma generators IPG3 and IPG4 are described here. An axial optical access through the inner injection head enables investigations of the plasma inside the generator. The tube cooling system is transparent. Hence, the position of the "plasma flame" within the tube can be observed with regard to different operating parameters such as chamber pressure, working gas, mass flow and anode power. Additionally, this feature is supported by the axial optical window. The total length of IPG3 is about 0.35 m, its diameter is about 0.1 m. The quartz tube contains the generated plasma which leaves the generator through the chamber adapter. The induction coil is connected to the external resonant circuit delivering power and cooling water for the IPGs. Furthermore, both the tube and the coil are surrounded by the tube cooling, which protects the quartz tube from overheating. The water and an additional cage around the generator serve as an rf-radiation protection shield. The length of IPG4 is about 0.4 m (with nozzle). IPG5 is an advanced design where reducing the distance between plasma and coil in turn further reduces the coupling losses as realized by Nawaz, Herdrich (2009).

2.2 Measurement Techniques

2.2.1 Cavity Calorimeter

A cavity calorimeter has been developed in order to determine the thermal plasma power in the PWK3 vacuum chamber (Fig. 5). The working principle is easy. The plasma enters the copper cavity, which is shaped like a cone, through a hole. The diameter of the hole is 120 mm, which is roughly 25 % bigger than the usual plasma beams. The distance between the calorimeter and the plasma outlet of the IPG is about 100 mm. This is necessary because smaller distances can result in retroactions that manipulate the discharge behaviour of the IPG; see

Nawaz, Herdrich (2009). Copper is used due to its very high specific heat conductivity and to its high catalysis. The copper walls are heated up through radiation, convection and recombination. The cavity is equipped with spiral copper tubes that guide the water that cools the copper wall.

With the measured cooling water exit temperature, the cooling water inlet temperature and the cooling water flow rate the plasma power can be determined:

$$P_{cal} = \rho_w c_{p,w} \dot{V}_w (T_{out} - T_{in}) + \Delta P \cdot \quad (5)$$

Here ρ_w is the density of the cooling water, \dot{V}_w is the cooling water volume flow rate, $c_{p,w}$ is the heat capacity, T_{out} is the cooling water outlet temperature and T_{in} is the cooling water inlet temperature. The parameter ΔP characterizes power losses due to the temperatures and velocities of the hot gas leaving the calorimeter's outlet. The corresponding kinetic power can be estimated using equilibrium models for the enthalpy together with the measured gas temperatures and mass flow rates, see Herdrich (2004).

A mean specific enthalpy can be calculated:

$$\bar{h} = \frac{P_{cal}}{\dot{m}} = \eta_{tot} \frac{P_A}{\dot{m}} \quad (6)$$

2.2.2 Steady State Heat Flux / Pitot Pressure Probe

Figure 6 shows a drawing of the steady state heat flux and Pitot pressure double probe that was used. The steady state heat flux is measured on an insert which can easily be changed. The heat flux sensor insert has a diameter of 26.5 mm according to the European standard geometry for heat shield material samples. Usually, copper inserts are used as a reference. For the steady state case, the flow rate through the insert \dot{V} and the temperature difference of the incoming and outgoing water ($T_{out}-T_{in}$) are measured, the latter by electrically insulated and shielded resistance thermometers Pt100. The heat flux per unit area is then given by

$$\dot{q} = \frac{c_{p,w} \rho_w \dot{V}_w (T_{out} - T_{in})}{A} \quad (7)$$

where $c_{p,w}$ is the heat capacity of water, ρ_w the water density and A the surface area of the probe exposed to the plasma. On the right side of Fig. 6 the Pitot pressure hole can be seen. For the present investigations the same diameter as for the heat flux insert was used. However, a variation of diameter can be performed by exchanging the modular insert for the pressure hole.

The Pitot pressure enables statements on the flow condition, in particular on the Mach number via the Rayleigh-Pitot relation, see e.g. Herdrich, Petkow (2008).

2.2 Processing of data to obtain mass specific enthalpies

On basis of known fully catalytic heat flux \dot{q}_{fc} and Pitot pressure the semi-empirical equation of Marvin and Pope, see Löhle, Lein, Eichhorn, Herdrich, Winter (2010)

$$h = \frac{\dot{q}_{fc}}{K} \left(\frac{p_{tot}}{R_{eff}} \right)^{-0.5} \quad (8)$$

can be used to calculate the mass specific enthalpy. Here, K is a gas specific constant and R_{eff} is an effective nose radius accounting for bluntness of the body contour as e.g. used for atmospheric entry. This effective radius depends on the incident flow condition and for low Mach numbers $R_{eff} \approx 2.3 \cdot R_{probe}$. If in Eq. (8) the dependency of both heat flux and pressure on the location was introduced with x as axial distance from the plasma generator outlet and y as radial position of the probe e.g. at a given, then the Integration of the enthalpy over dA with an assumed constant mass flow density yields the plasma power P_{Pl} :

$$P_{Pl} = \sqrt{R_{eff}} \frac{2\pi \dot{m}}{K A} \int_{y=0}^{y=R_{Pl}} \frac{\dot{q}_{fc}(y)}{\sqrt{p_{tot}(y)}} y dy \quad (9)$$

Eq. (9) can be used to replace K in Eq. (8) which leads to:

$$\frac{h(x, y)}{h_{eff}} = \frac{R_{pl}^2}{2} \frac{\dot{q}_{fc}(x, y)}{\sqrt{P_{tot}(x, y)}} \int_{y=0}^{y=R_{pl}} \frac{\dot{q}_{fc}(y)}{\sqrt{P_{tot}(y)}} y dy \quad (10)$$

3. Review of CO₂ Plasma Conditions for PWK3

An extensive database for experimental conditions of both CO₂ and CO₂ with N₂ has been developed in the past at IRS using the above described plasma wind tunnel PWK3. In addition, experiments have been conducted with iron oxide powder in order to simulate the dusty atmosphere of Mars by Endlich (2008). These data are not included in this analysis. Further measurements described this reference imply a significant influence of the thermochemical situation as the heat fluxes measured with an iron insert were about 25% larger than the values resulting from the copper base. However, the analysis showed that this additional heat flux cannot be necessarily traced back to catalysis only and that additional exothermal reactions such as e.g. oxidation may took place.

Data obtained by Herdrich are used to calculate the empiric constant K by Marvin and Pope. Having this constant the heat flow of large variety of experiments can be transformed to mass specific enthalpies. Therefore, a comparison of the conditions during Mars and Venus entry and those being generated by the plasma generators at IRS can be conducted.

3.1 Gas specific K from conditions with known enthalpies

Using the cavity calorimeter the plasma power has been measured and hence, knowing the mass flow, the effective specific enthalpy. Additionally, the radial heat flux and pressure profiles have been measured at five axial distances to the plasma generator outlet. Thus, the specific enthalpies could be calculated using the relation given by Eq. (9). Regarding the determined specific enthalpies, heat fluxes and total pressures on the centreline of the generator the constant K by Marvin and Pope has been determined by a least square approach, when calculating the enthalpy with Eq. (7). The resulting value is shown in Table 1 compared with two values derived from algebraic boundary layer calculations given by Löhle, Lein, Eichhorn, Herdrich, Winter (2010). Additionally a corrected value implying the catalysis of copper oxide is presented. The fully-catalytic heat flux is about 10 % higher than the heat flux on copper oxide, see Kolesnikov, Yakushin, Pershin, Vasil'evskii, Bykova, Gordeev (2002).

The experimentally determined values for K are about 19 % lower than the calculated values given by Löhle, Lein, Eichhorn, Herdrich, Winter (2010). However, these values have been determined using boundary layer equations based on the derivations of Fay and Riddell. Since this equation takes not well known parameters into account which are often approximated by equilibrium relations (as for example the viscosities at the wall and the boundary layer edge) uncertainties evolve that may explain the resulting differences. However, the correspondence between the two methods is fairly meaning that the overall results verify each other. This by the way eventually motivated Marvin and Pope to develop their simplified semi-empirical relation as the overall applicability of the aforementioned boundary layer equations were accompanied by such sensitivities. The enthalpies presented throughout this paper have been calculated using the equation by Marvin and Pope using the K derived by the experimental data presented within this paper.

3.2 Derivation of mass specific enthalpies for further plasma conditions

For the majority of the existing experimental data this K has been used to estimatively provide enthalpies for the displayed conditions (Eq. (8)). The overall data base is then compared with relevant conditions during Mars and Venus entries. The data has been used to characterize the feasibility of CO₂ plasma conditions relevant for Mars or Venus entry maneuvers using PWK3. In Fig. 7 a comparison of the measured data within this paper and known trajectories is presented.

4. Operational Envelopes of PWK3 for Mars and Venus

The data of multiple operating points has been used to find the operational range of PWK3. The thermochemical simulation requires the duplication of the three key parameters total enthalpy, stagnation pressure and velocity gradient as e.g. outlined by Kolesnikov (2000) and Kolesnikov, Marraffa (1999). As the similarity of the velocity gradient can be achieved by geometrical scaling of the probe, the pressure and enthalpy ranges define the operational range of the plasma wind tunnel. Figure 7 shows the results in comparison with those enthalpy and pressure ranges being of interest for entries into the atmospheres of Mars and Venus. On the x-axis both the specific enthalpy and the velocity during entry are shown. The velocity of the plasma flow is smaller. However,

the mass specific enthalpies are duplicated.

For Mars, the trajectories of the Phoenix and Pathfinder Lander and a typical ballistic trajectory are shown. For Venus the trajectories of the Pioneer Venus Day and Night Probes are shown, see Seiff, Kirk (1982). Additionally, lines of equal altitude of the Mars and Venus atmosphere are plotted. For Mars, the atmosphere is modeled isothermal using density and pressure data obtained by Spirit during descend at an altitude of 20 km. For Venus the obtained atmospheric entry data of the Pioneer Venus Day Probe has been used (see Seiff, Kirk (1982)). In terms of pressure the currently performed operational regime of PWK3 is between 1.3 hPa and 20 hPa. The maximum achieved mass specific enthalpy is about 63 MJ/kg. By decreasing the mass flow rates for the plasma generator even higher specific enthalpies would be possible. The minimum pressure is limited on the one hand by the vacuum pump performance and on the other hand by the plasma generator.

The operational range allows simulating the conditions during early Mars entry, including peak deceleration, matching both enthalpy and pressure conditions. In addition, the enthalpy conditions can be reproduced for both Mars and Venus entries at the whole. However, it is evident that the facility should be extended to a higher pressure regime in future as both filling the gap to the higher pressure regimes and the development of equilibrium and near equilibrium test cases e.g. for plasma radiation research can be performed.

All experimental data were taken from the following references: Endlich, P. (2008), Kuhn (2001), Herdrich, Auweter-Kurtz, Endlich, Laux, (2003) and Endlich, Auweter-Kurtz, Herdrich, Löhle, Winter, (2003). In addition, enthalpy conditions from other facilities, all of them inductively driven except the arc-driven L2K, are shown. The values used for this are reported in the following references: Koch, Riemer, Esser, Gülhan, (2009), Gülhan, Esser, Koch, (2002), Bykova, Vasilevskii, Gordeev, Kolesnikov, Pershin, Yakushin (1997), Kolesnikov, Pershin, Vasilevskii (2004), Vacher, Andre, Dudeck (2008) and Rond, Bultel, Boubert, Cheron (2008). With the L2K data an attempt to create a data base with maximum completeness is made although these facilities typically suffer from electrode problems that accompany the use of CO₂ as working gas.

5. Conclusions

Based on measurements of radial total pressure and heat flux profiles of CO₂ plasma in PWK3 the constant K introduced by Marvin and Pope could be determined experimentally. Comparison with values being derived on basis of the Fay-Riddell equation showed deviations of up to 38 %. The respective reasons for these deviations have been discussed. The obtained value for K has been used to determine the specific enthalpies of previously measured CO₂ plasma conditions. Those enthalpies have been compared with conditions during Mars and Venus entry. The comparison revealed that PWK3 is able to provide the enthalpy conditions of the whole Mars and Venus entry. The rather lower pressures for the atmospheric Mars entries allow simulating both pressures und specific enthalpies fairly well for this scenario. Venus entries are out of the pressure range. Consequently the extension of the pressure regime of PWK3 will be subject to further research.

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Table 1. Comparison of different values for K

K _{CO₂} [kW/(m ^{3/2} Pa ^{1/2} MJ)]	Comments
0.28	copper based heat flux
0.31	factor of 1.1 used to obtain fully-catalytic heat flux, see text
0.37	Ref. [16] based on Fay-Riddell
0.43	Ref. [19] based on Fay-Riddell

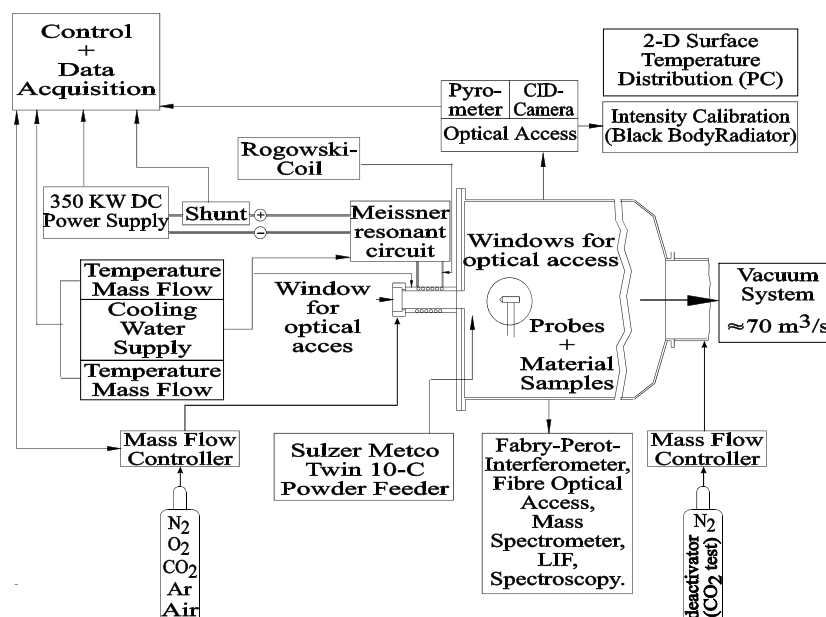


Figure 1. PWK3 facility set-up

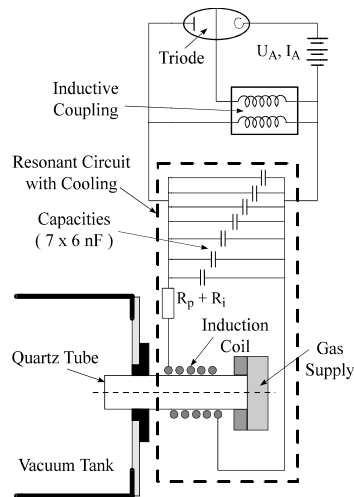


Figure 2. Scheme of Meissner type resonant circuit

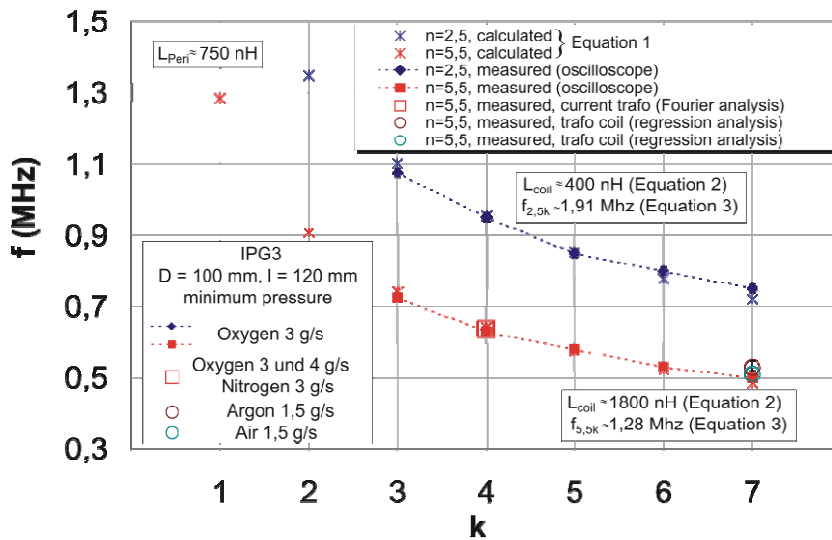


Figure 3. Measured operational frequencies ($n = 2.5, n = 5.5$) depending on number of capacitors

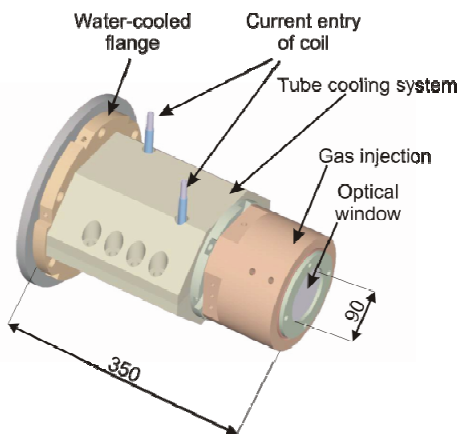


Figure 4. Inductively heated plasma source IPG3: assembly (left), photograph of PWK3 (right)

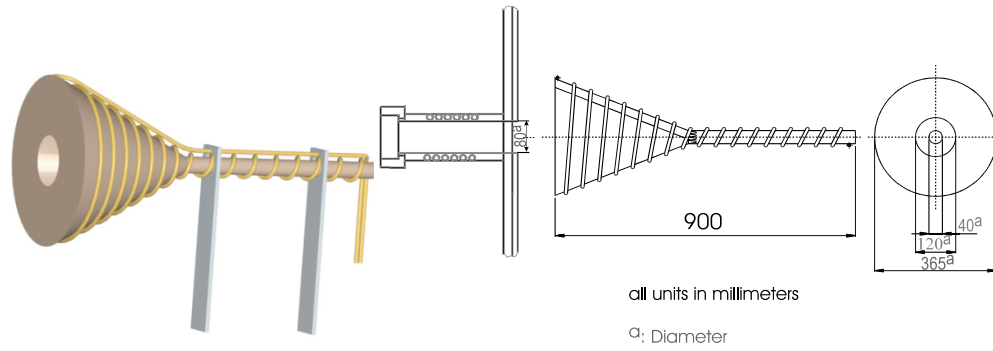


Figure 5. Calorimeter: 3-dimensional drawing (left), scheme in PWK3 with dimensions (right)

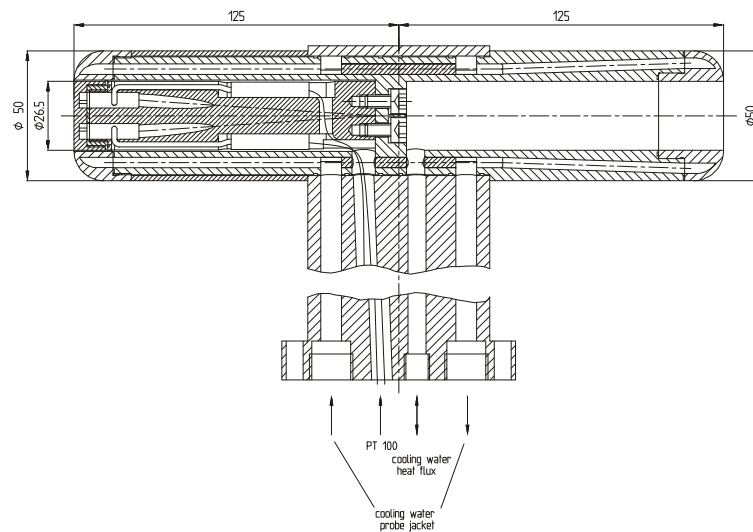


Figure 6. Scheme of Steady State Heat Flux and Pitot Pressure Double Probe

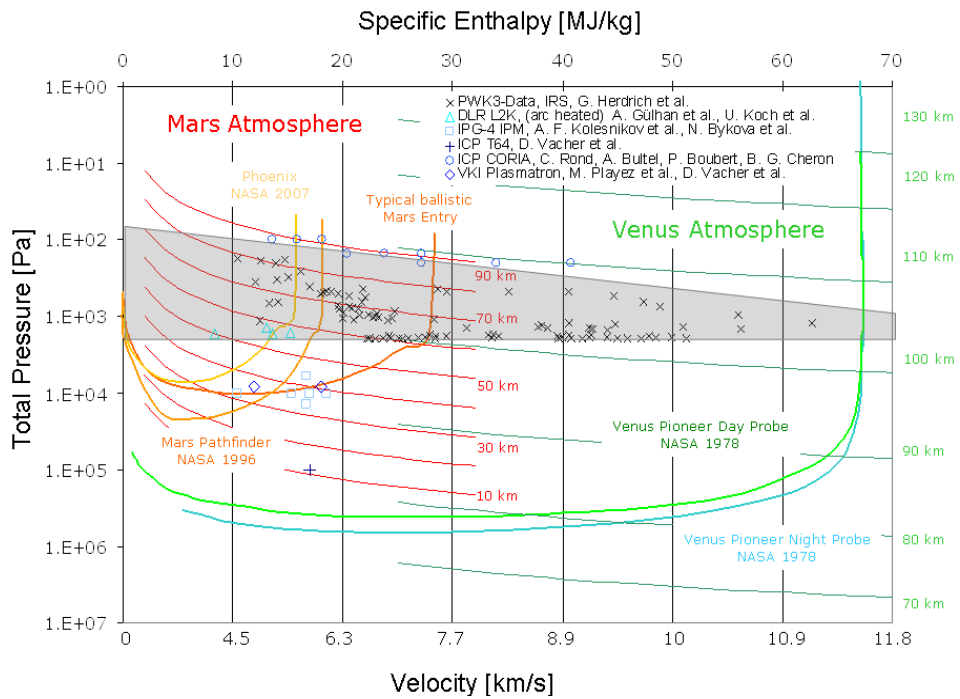


Figure 7. Operational conditions of PWK3 compared with Mars and Venus entry trajectories and enthalpy data for other facilities